

# 非一様物質下での ニュートリノ振動におけるパラメータ励振

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## Plan

1. Introduction
2. Formulation
3. Parametric Resonance in Oscillation Probability
4. Summary & Outlook

# Introduction

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# Neutrino Oscillation

## • Discovery, establishment of the oscillation

- Solar Homestake, Kamiokande...; Super-K, SNO, ...
- Atmospheric Kamiokande, ...; Super-K, ...
- Reactor Chooz, KamLAND,...; Double Chooz,...
- Accelerator K2K, MINOS; T2K,...

$$\delta m_{\text{atm}}^2 = 2.39 \left( 1_{-0.084}^{+0.013} \right) \times 10^{-3} \text{eV}^2 \quad (2\sigma) \quad \sin^2 \theta_{12} = 0.312 \left( 1_{-0.109}^{+0.128} \right) \quad (2\sigma)$$

$$\delta m_{\text{sol}}^2 = 7.67 \left( 1_{-0.047}^{+0.044} \right) \times 10^{-5} \text{eV}^2 \quad (2\sigma) \quad \sin^2 \theta_{23} = 0.466 \left( 1_{-0.215}^{+0.292} \right) \quad (2\sigma)$$

$$\delta_{\text{CP}} \quad \sin^2 \theta_{13} = 0.012 \pm 0.013 \quad (1\sigma)$$

Fogli et al. (2008)

# Neutrino Oscillation

## • Precision Measurement

- Tiny unknown parameters:  $\theta_{13}$ ,  $\delta_{CP}$
- Pin-down the parameter values
- Lifting degeneracy

## • Reactor

## • Accelerator + Very-long-baseline

- $L \gtrsim 1000\text{km} + 2$  Detectors
- $L \gtrsim 2000\text{km}$
- $L \gtrsim 7000\text{km}$
- etc.

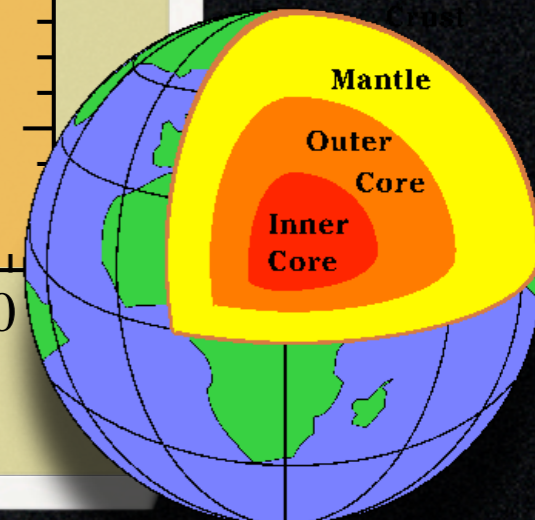
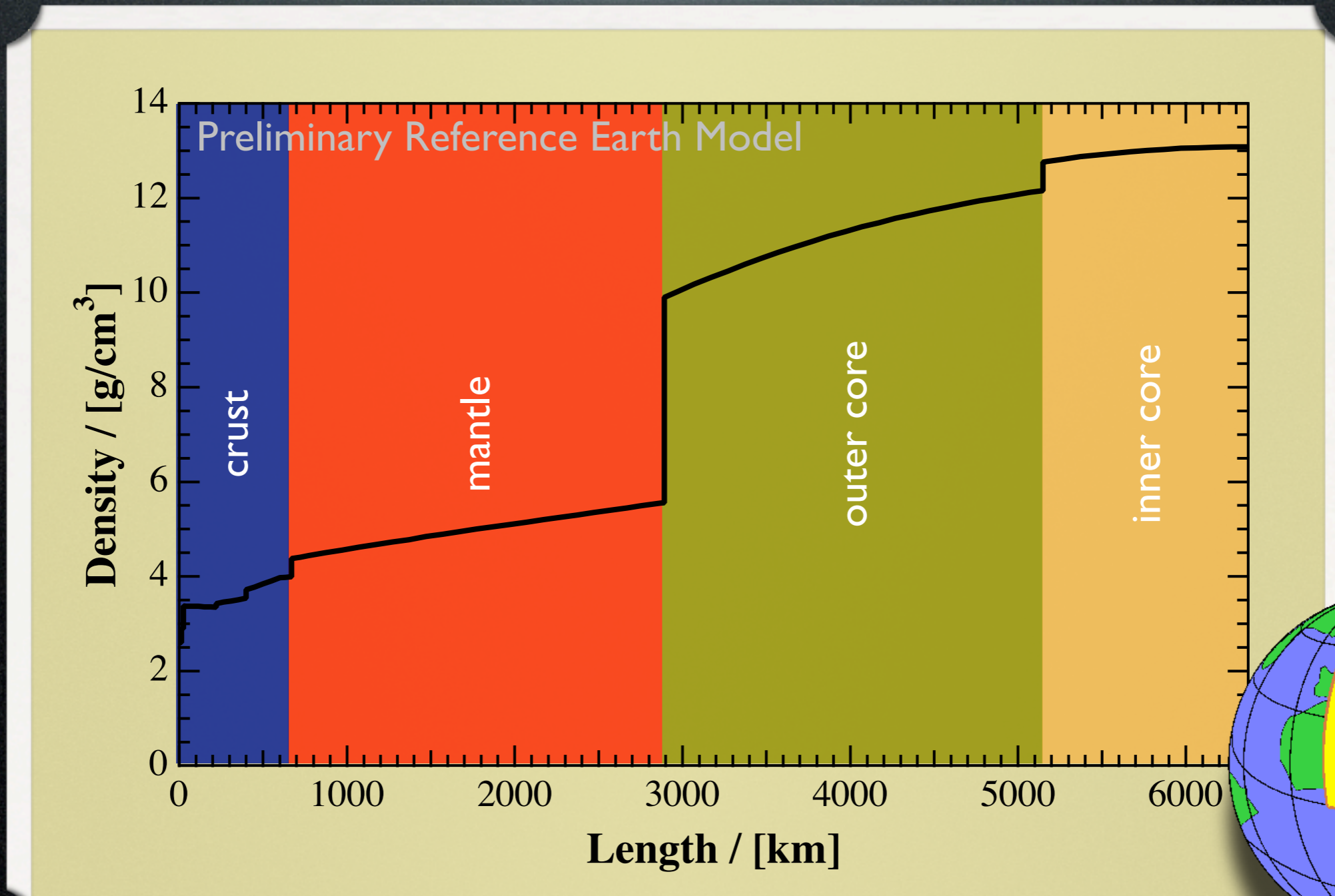
—————→ Matter effect in control

## • Others

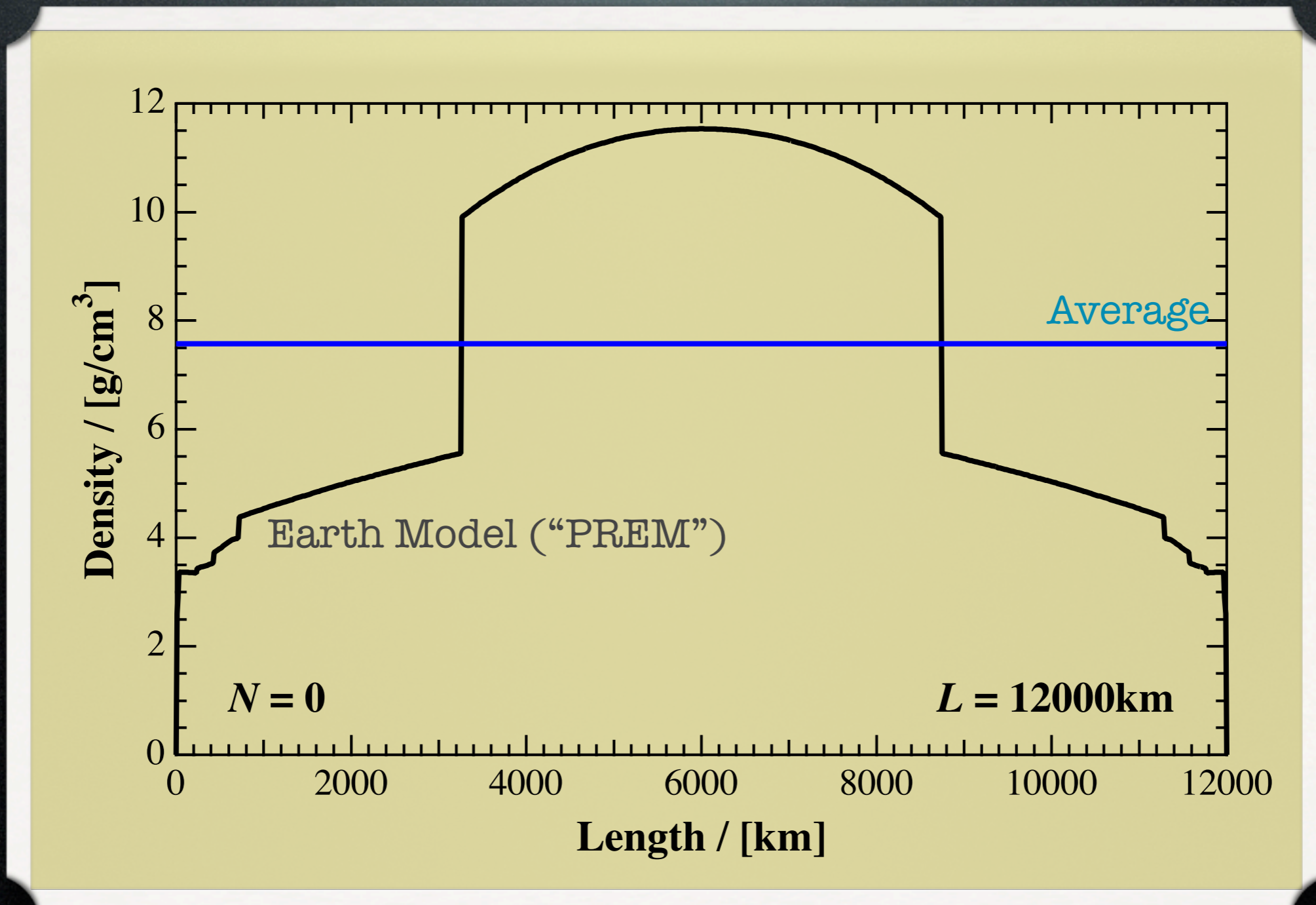
# Matter effect

- Matter density on the baseline
  - Constant
  - Tabulation according to an Earth model
  - Analytic approximation

# Earth Model, Example

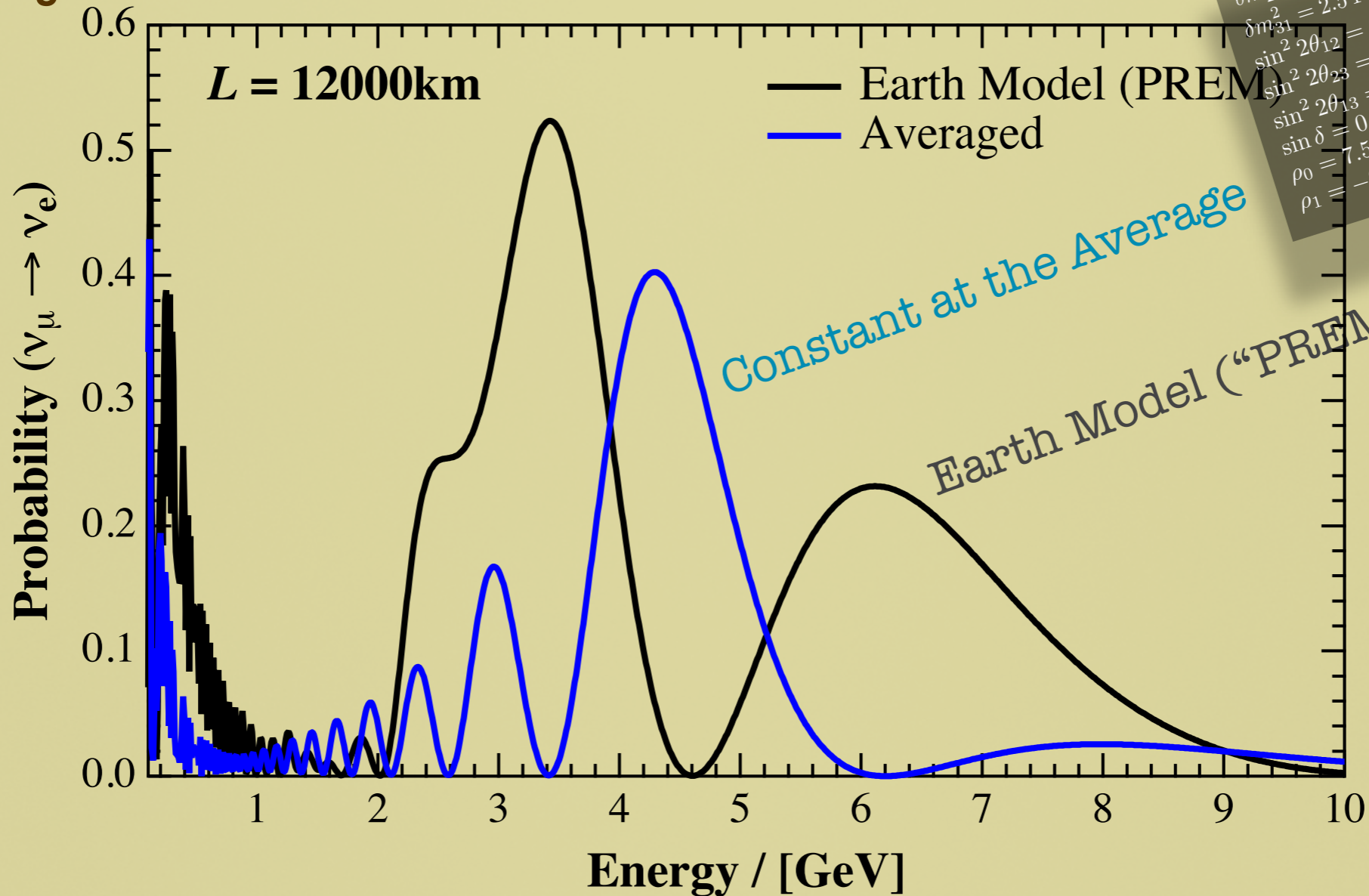


# Matter Density on a Baseline



# Constant vs. Earth Model

## Qualitative difference?



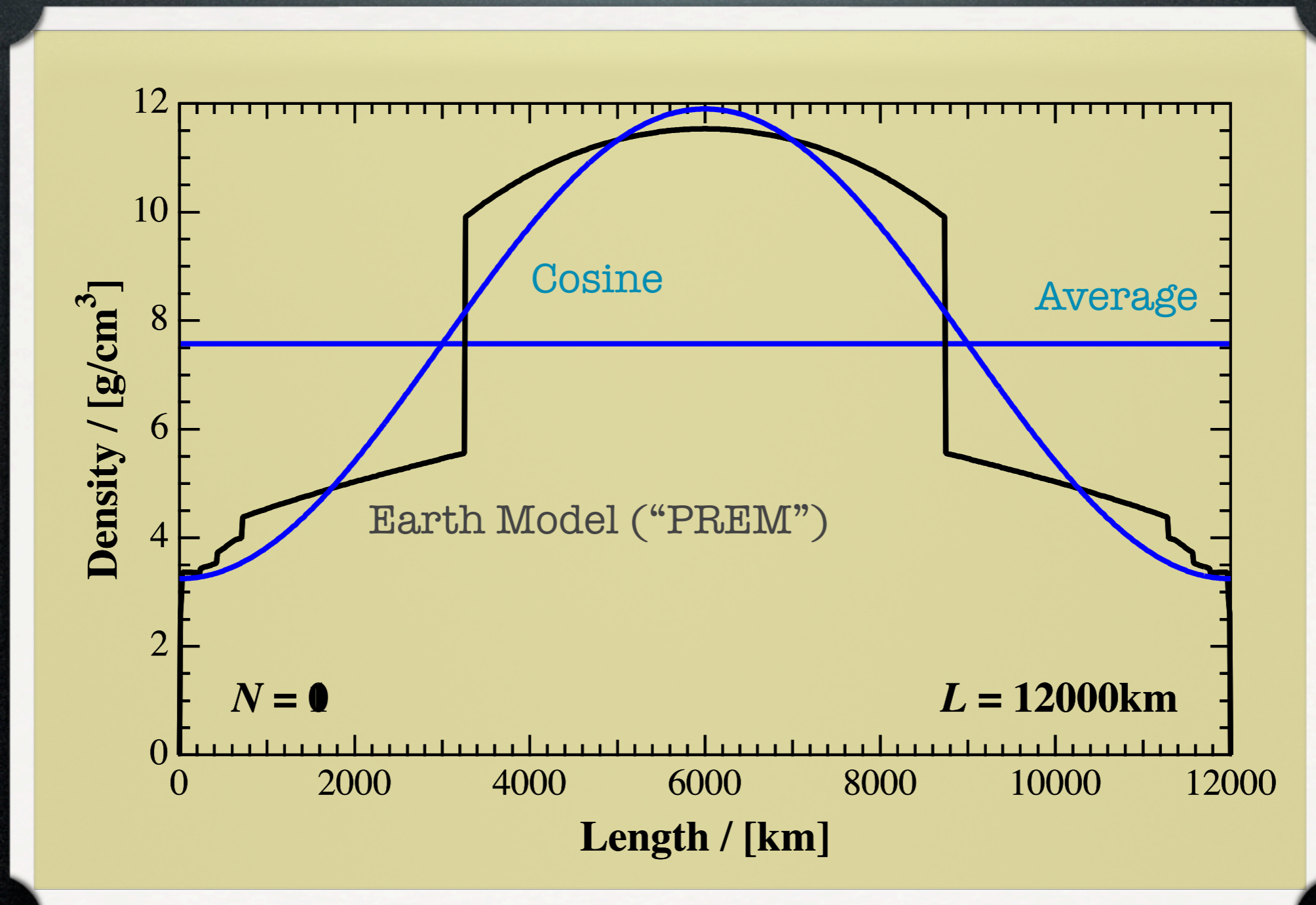
$L = 12000\text{ km}$   
 $\delta m_{21}^2 = 7.9 \cdot 10^{-5} \text{ eV}^2$   
 $\delta m_{31}^2 = 2.5 \cdot 10^{-3} \text{ eV}^2$   
 $\sin^2 2\theta_{12} = 0.84$   
 $\sin^2 2\theta_{23} = 1.00$   
 $\sin^2 2\theta_{13} = 0.05$   
 $\sin \delta = 0.00$   
 $\rho_0 = 7.58 \text{ g/cm}^3$   
 $\rho_1 = -2.16 \text{ g/cm}^3$

Constant at the Average

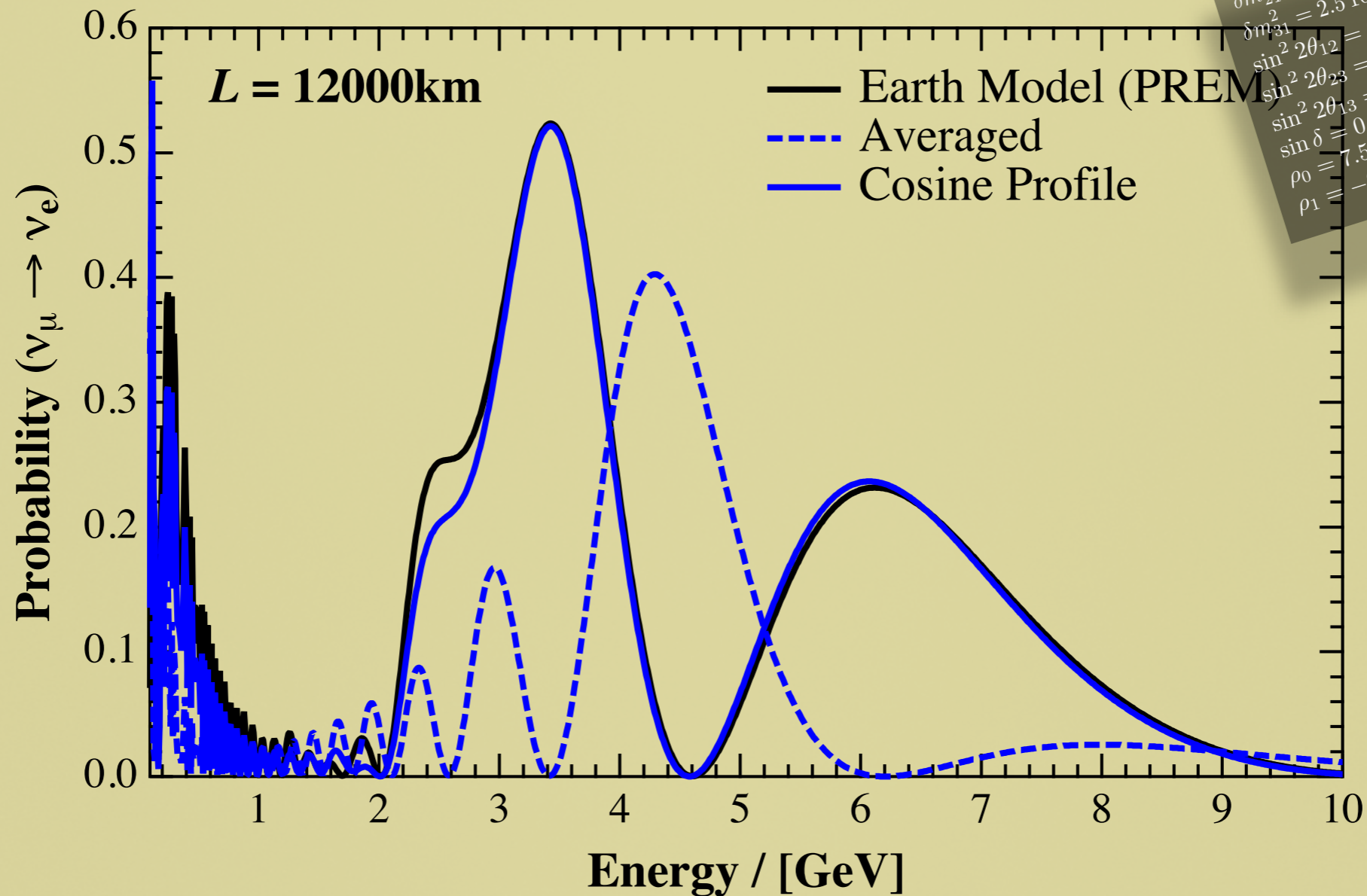
Earth Model ("PREM")



# Matter Density Profile

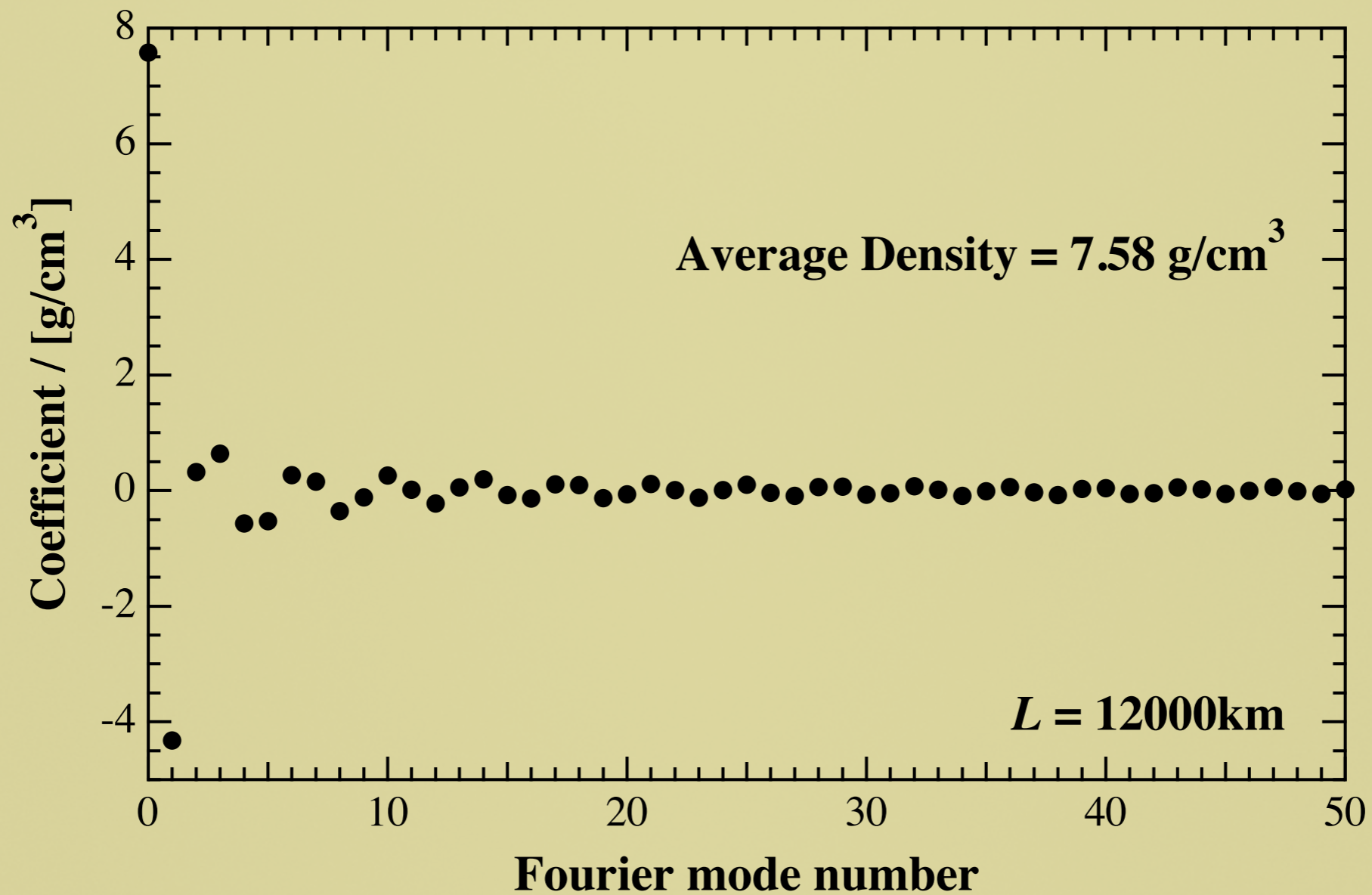


# Constant vs. Earth Model

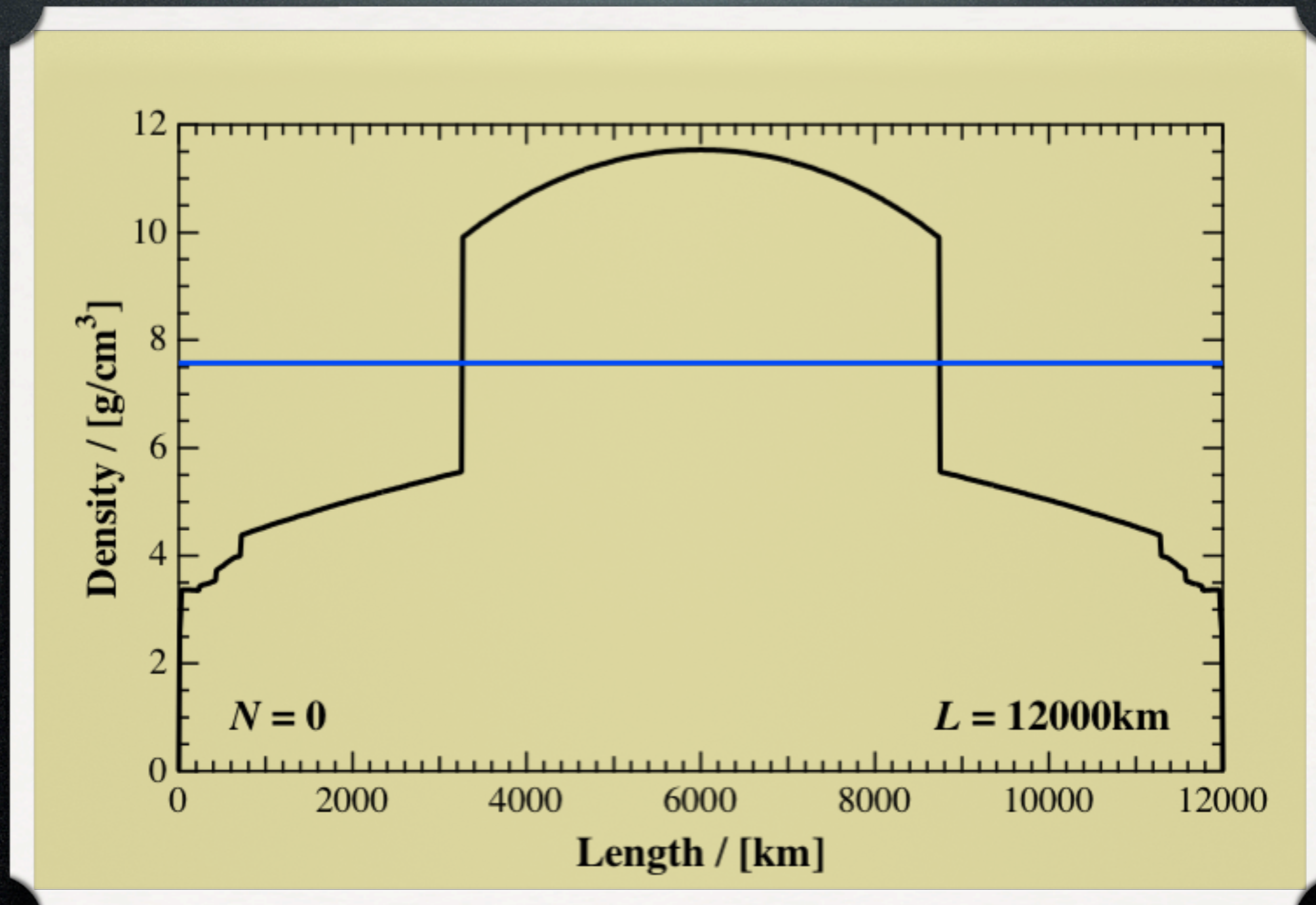


$L = 12000\text{ km}$   
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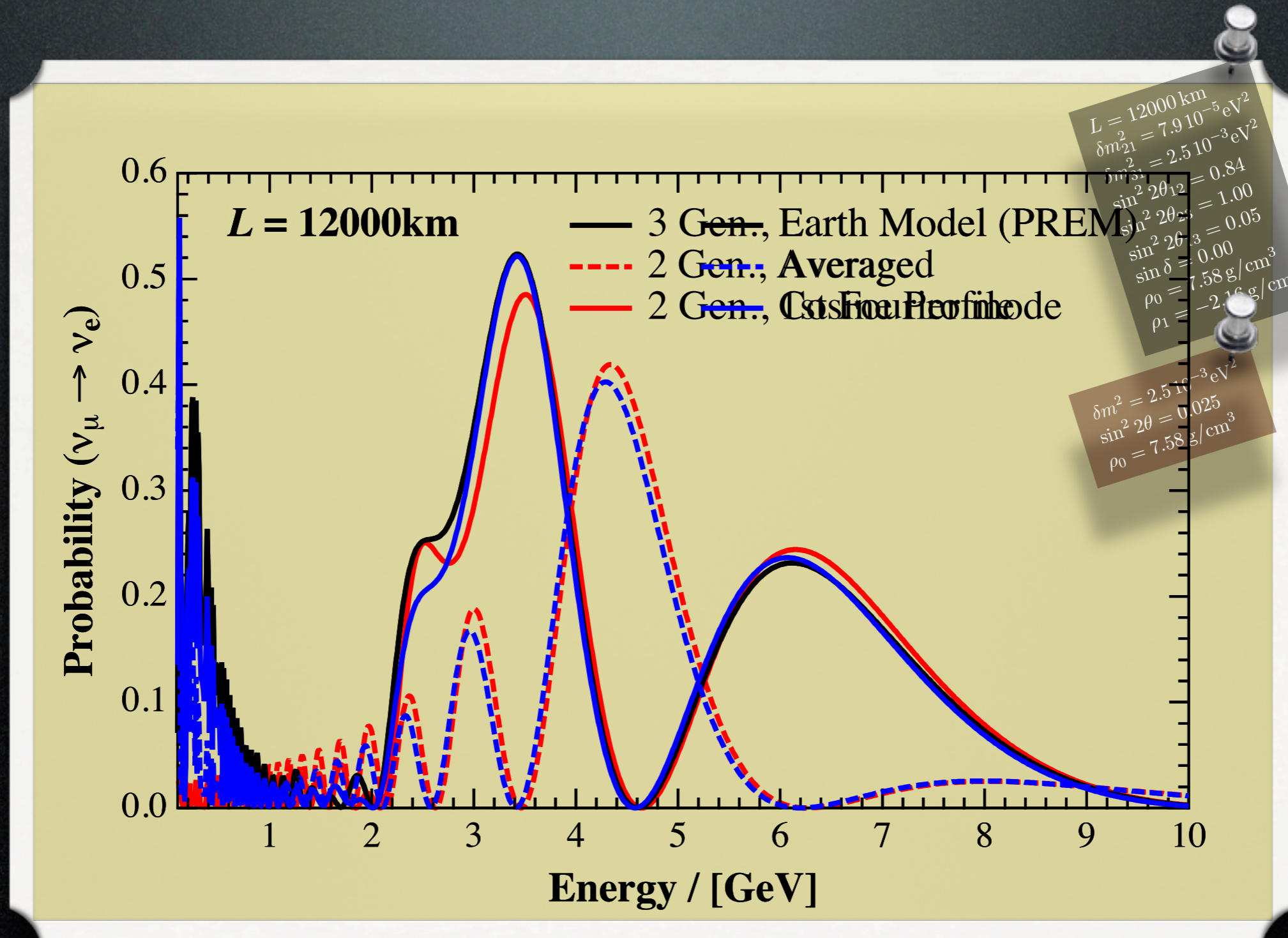
# Fourier Coefficients



# Matter Density Profile



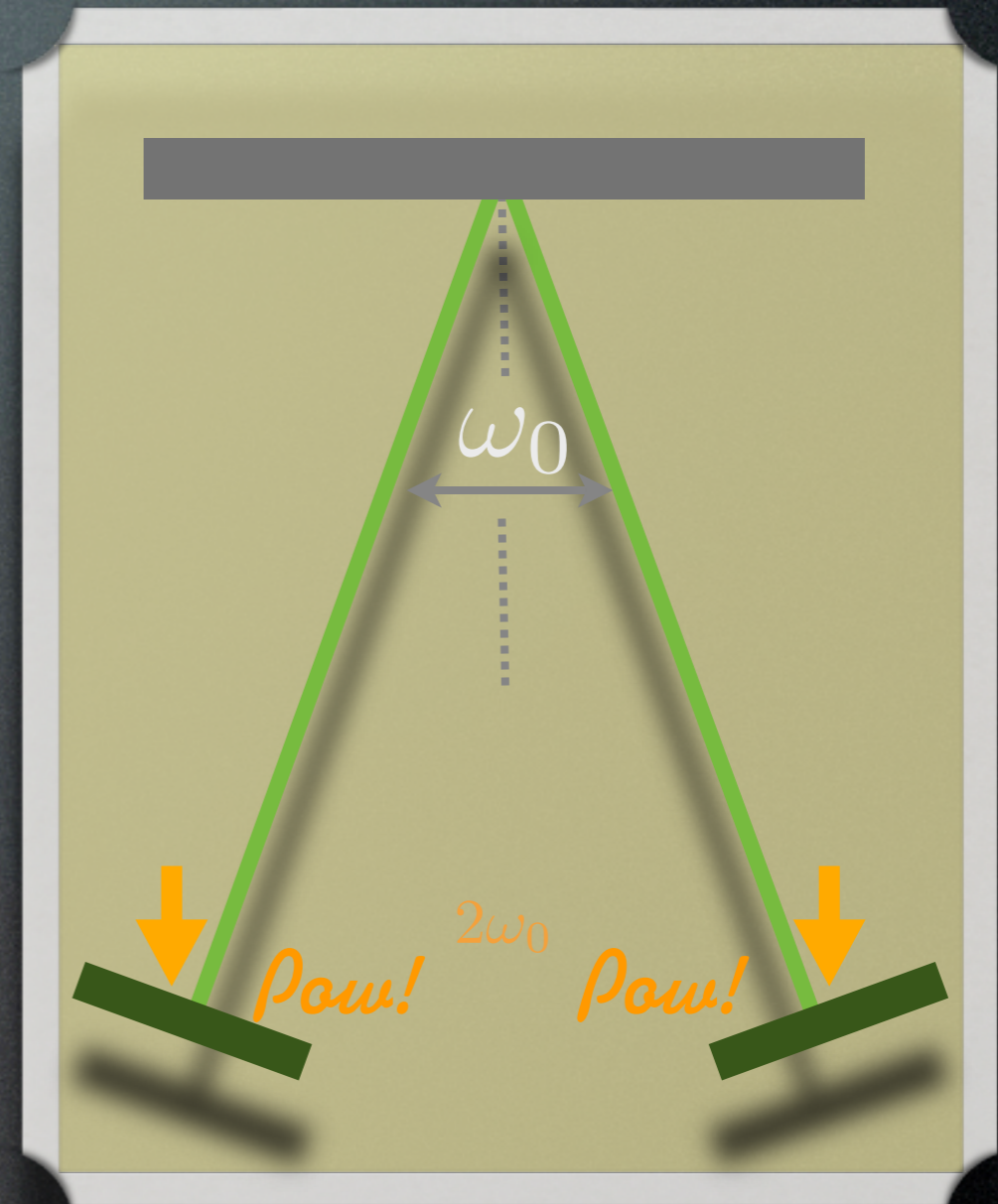
# We'll proceed with 2 flavors



# Parametric Resonance

- Periodic perturbation
  - Twice in a period
  - Grows amplitude of oscillation
- Matter effect as a bunch of periodic perturbations

Ermilova et al. (1986), Akhmedov (1988), Krastev&Smirnov (1989), Krastev&Smirnov (1989), Liu&Smirnov (1998), Petcov (1998), Chizhov&Petcov (1998), ..., Akhmedov&Maltoni&Smirnov (2005), ...

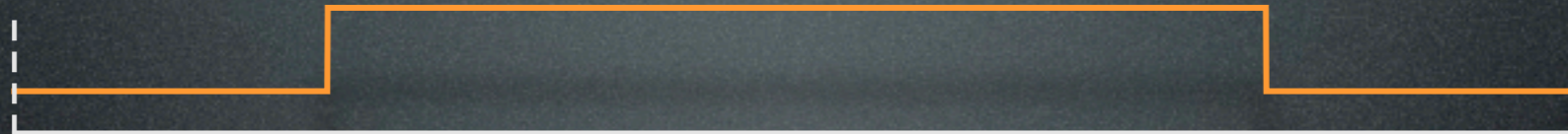


# Formulation



# Modeling Density Profiles

## • Step function



Akhmedov (1988), Krastev&Smirnov (1989), Krastev&Smirnov (1989),  
Liu&Smirnov (1998), Petcov (1998), Chizhov&Petcov (1998), ...,  
Akhmedov&Maltoni&Smirnov (2005), ...

## • Fourier series



- Koike&Sato (1998), Ota&Sato (2003), ...



# Two-Flavor Oscillation

- Evolution equation of the two-flavor neutrino

$$i \frac{d}{dx} \begin{pmatrix} \nu_e(x) \\ \nu_\mu(x) \end{pmatrix} = \frac{1}{2E} \left[ \frac{\delta m^2}{2} \begin{pmatrix} -\cos 2\theta & \sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} + \begin{pmatrix} a(x) & 0 \\ 0 & 0 \end{pmatrix} \right] \begin{pmatrix} \nu_e(x) \\ \nu_\mu(x) \end{pmatrix}$$

- Matter effect  $a(x) = 2\sqrt{2}G_F n_e(x)E$

- Dimensionless variables:

$$\begin{array}{ccc} \xi \equiv \frac{x}{L} & \Delta \equiv \frac{\delta m^2 L}{2E} & \Delta_m(\xi) \equiv \frac{a(\xi)L}{2E} \\ \text{Distance} & \text{Reciprocal } E & \text{Matter Density} \end{array}$$

- $z(\xi) = \nu_e(\xi) \exp \left[ \frac{i}{2} \int_0^\xi ds \Delta_m(s) \right] \quad \dots \quad |\nu_e(\xi)|^2 = |z(\xi)|^2$

- Initial conditions  $\nu_e(0) = 0, \nu_\mu(0) = 1 \rightarrow z(0) = 0, z'(0) = -i \frac{\Delta}{2} \sin 2\theta$

- Second-order equation in dimensionless variables

$$z''(\xi) + \frac{1}{4} \left[ (\Delta_m(\xi) - \Delta \cos 2\theta)^2 + \Delta^2 \sin^2 2\theta + 2i\Delta'_m(\xi) \right] z(\xi) = 0$$

# Constant-Density Matter

$$z''(\xi) + \frac{1}{4} \left[ (\Delta_m(\xi) - \Delta \cos 2\theta)^2 + \Delta^2 \sin^2 2\theta + 2i\Delta'_m(\xi) \right] z(\xi) = 0$$

- Simple solution when  $\Delta_m(\xi) \equiv \Delta_0 = (\text{const.})$

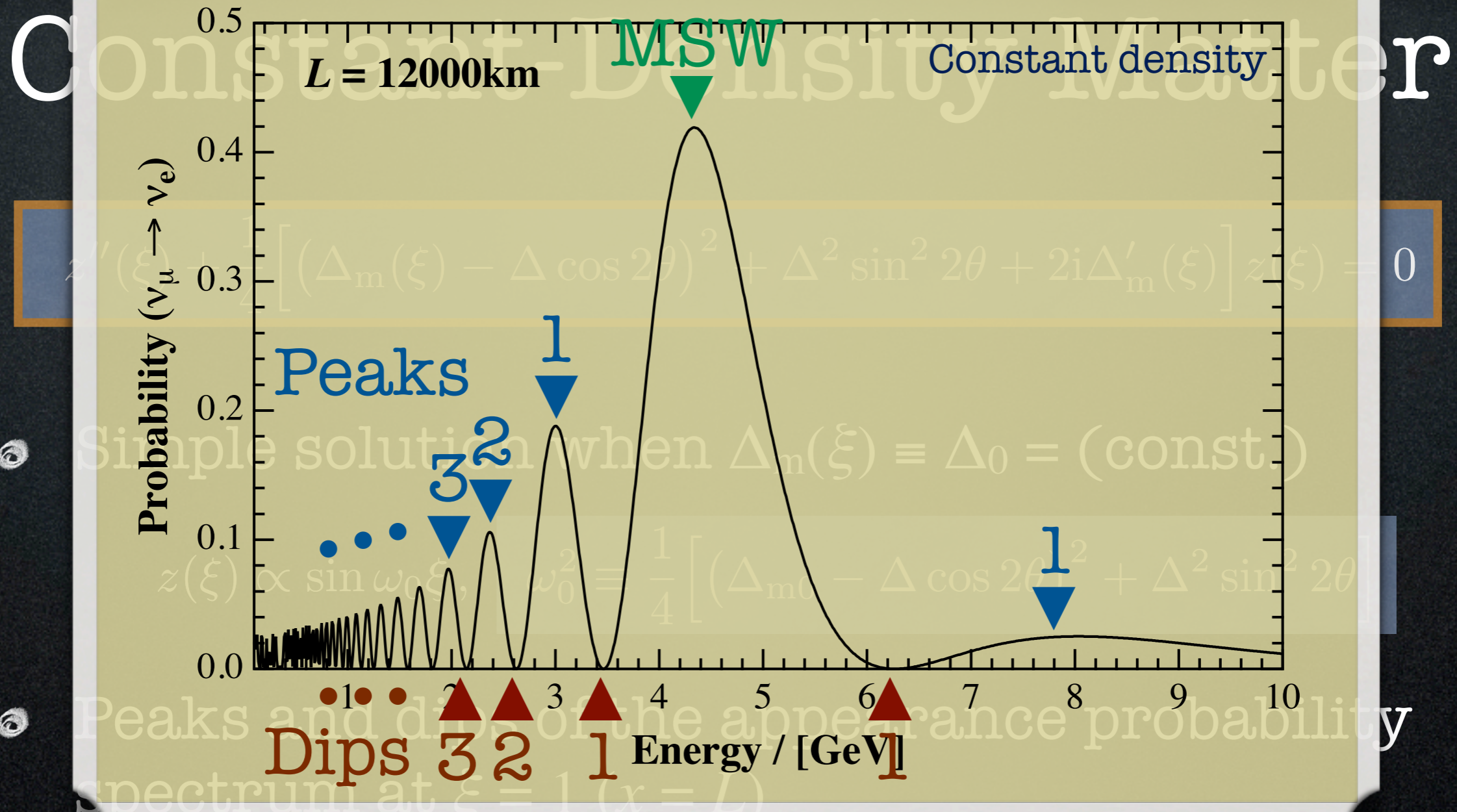
$$z(\xi) \propto \sin \omega_0 \xi, \quad \omega_0^2 \equiv \frac{1}{4} \left[ (\Delta_{m0} - \Delta \cos 2\theta)^2 + \Delta^2 \sin^2 2\theta \right]$$

- Peaks and dips of the appearance probability spectrum at  $\xi = 1$  ( $x = L$ )

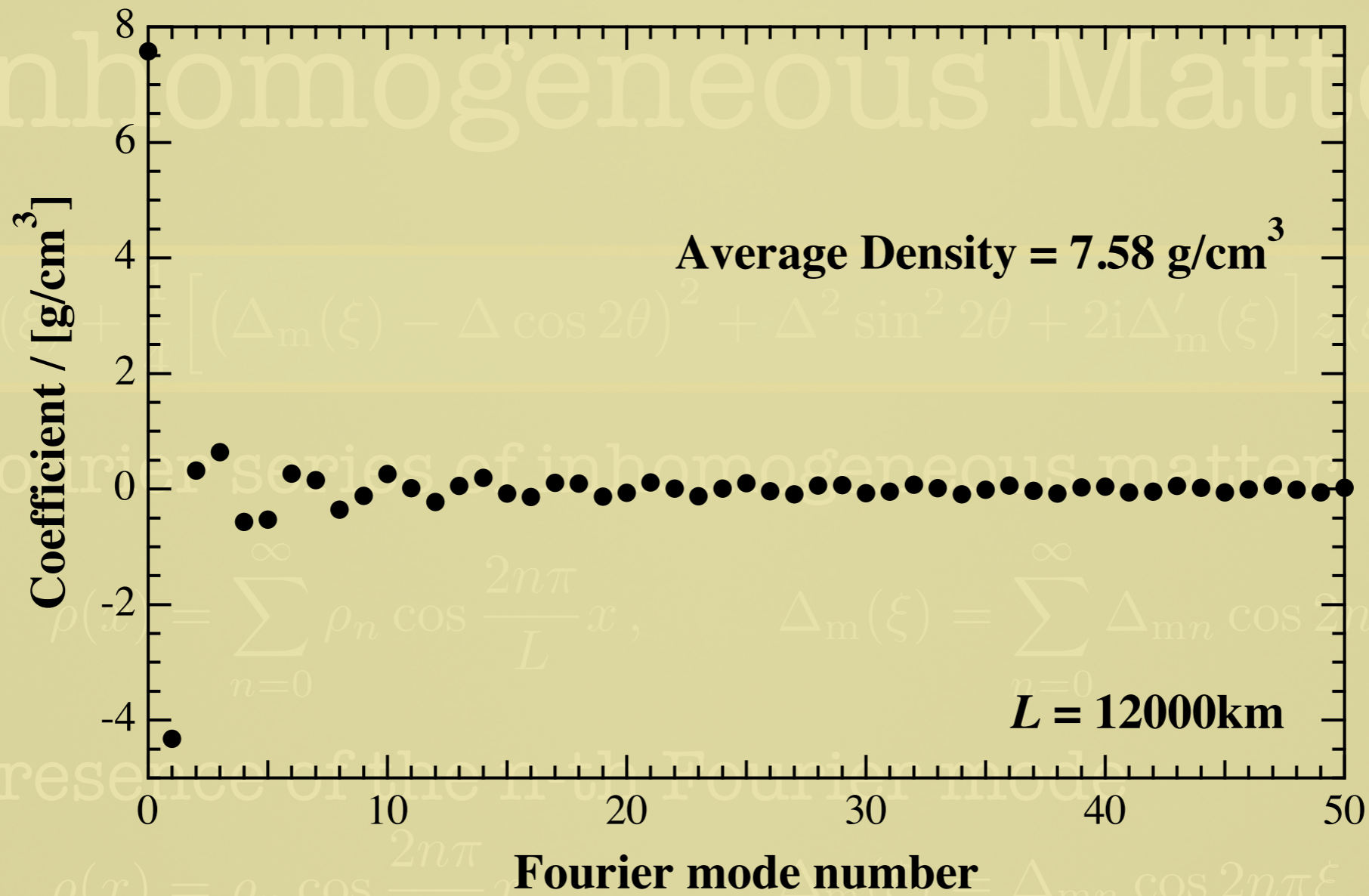
- $\omega_0 = \frac{1}{2} \Delta \sin 2\theta$  MSW resonance.

- $\omega_0 = \left( n + \frac{1}{2} \right) \pi$   $(n+1)$ -th peak.

- $\omega_0 = n\pi$   $n$ -th dip.



- $\omega_0 = \frac{1}{2} \Delta \sin 2\theta$       MSW resonance.
- $\omega_0 = \left(n + \frac{1}{2}\right) \pi$       (n+1)-th peak.
- $\omega_0 = n\pi$       n-th dip.



0

# Inhomogeneous Matter

$$z''(\xi) + \frac{1}{4} \left[ (\Delta_m(\xi) - \Delta \cos 2\theta)^2 + \Delta^2 \sin^2 2\theta + 2i\Delta'_m(\xi) \right] z(\xi) = 0$$

- Fourier series of inhomogeneous matter

$$\rho(x) = \sum_{n=0}^{\infty} \rho_n \cos \frac{2n\pi}{L} x, \quad \Delta_m(\xi) = \sum_{n=0}^{\infty} \Delta_{mn} \cos 2n\pi\xi$$

- Presence of the  $n$ -th Fourier mode

## Mathieu Equation

$$z''(t) + (\omega^2 - 2\varepsilon \cos t) z(t) = 0$$

$$z''(\xi) + (\omega_0^2 + \alpha_n \cos 2n\pi\xi - i\beta_n \sin 2n\pi\xi + \gamma_n \cos 4n\pi\xi) z(\xi) = 0$$

$$\omega_0^2 = \frac{1}{4} (\Delta_{m0} - \Delta \cos 2\theta)^2 + \frac{1}{4} \Delta^2 \sin^2 2\theta + \frac{1}{8} \Delta_{mn}^2,$$

$$\alpha_n = \frac{1}{2} (\Delta_{m0} - \Delta \cos 2\theta) \Delta_{mn}, \quad \beta_n = n\pi \Delta_{mn}, \quad \gamma_n = \frac{1}{8} \Delta_{mn}^2$$

# Parametric Resonance in Appearance Probability

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# Condition of Resonance

$$z''(\xi) + (\omega_0^2 + \alpha_n \cos 2n\pi\xi - i\beta_n \sin 2n\pi\xi + \gamma_n \cos 4n\pi\xi)z(\xi) = 0$$

n-th mode

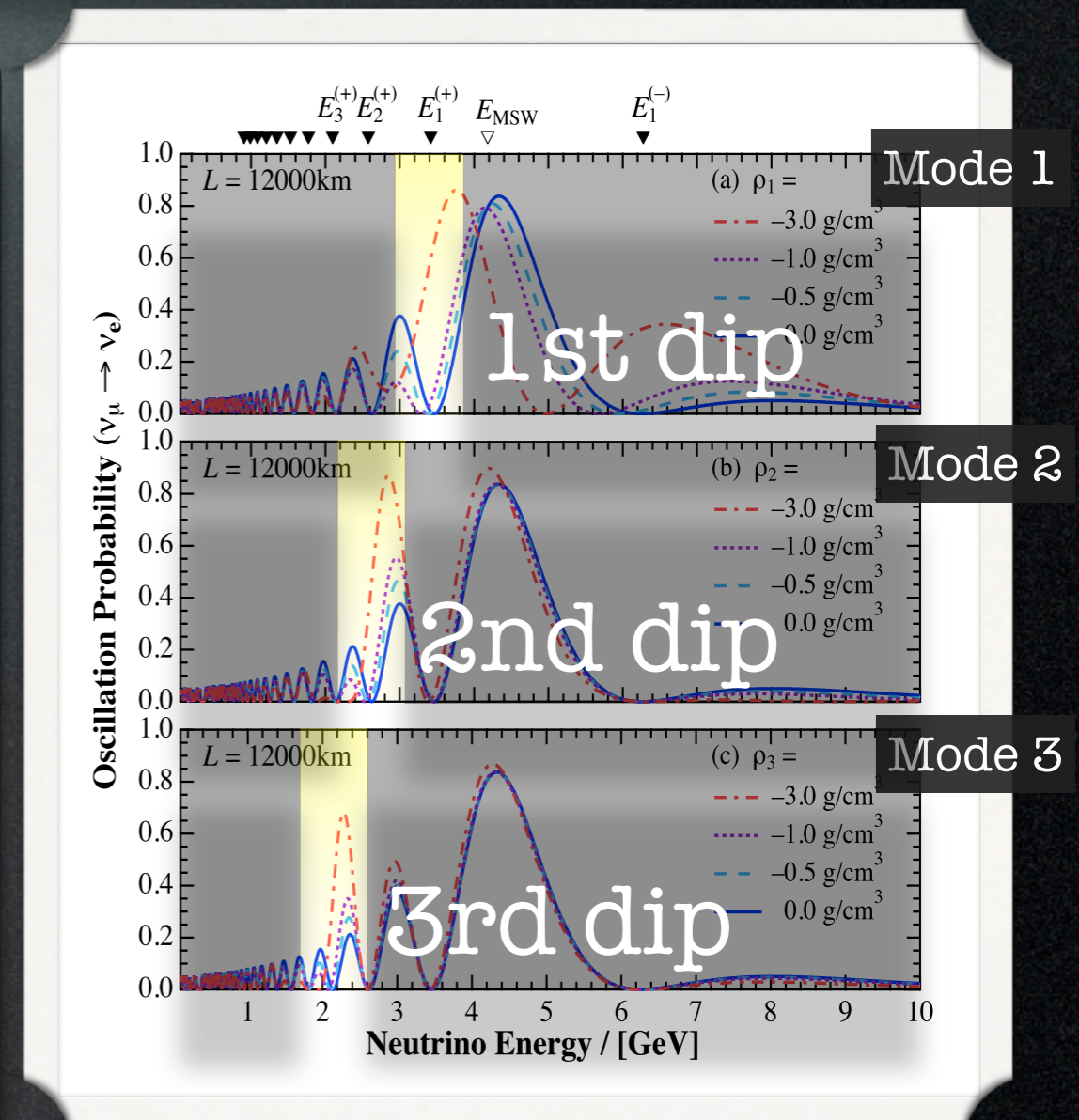
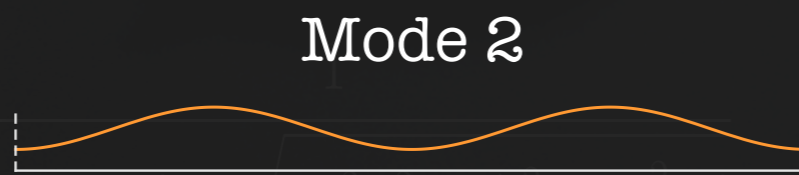
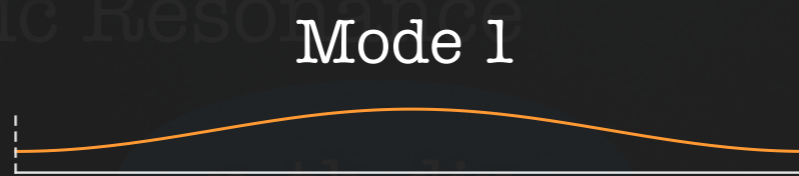
## Matter Profile

- Parametric Resonance Condition

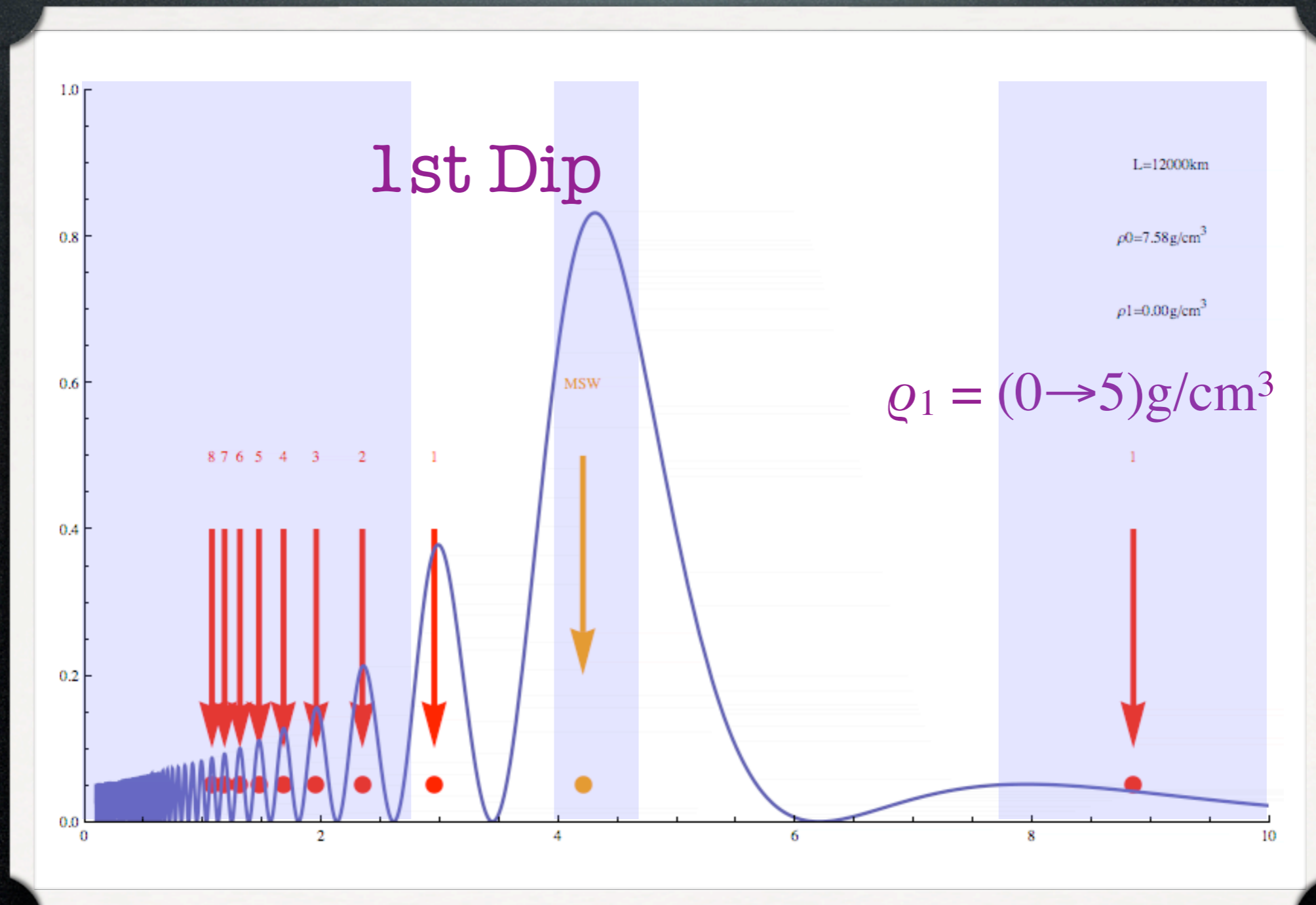
$$\omega_0 = n\pi$$

$$E = E_n^{(\pm)} \equiv \frac{\delta m^2 L}{2 \Delta_{1n} \cos 2\theta_{1n} \pm \sqrt{(\Delta_{1n} \sin 2\theta_{1n})^2 + (\delta m^2)^2}}$$

- n-th Fourier profile affects the appearance dip of the a

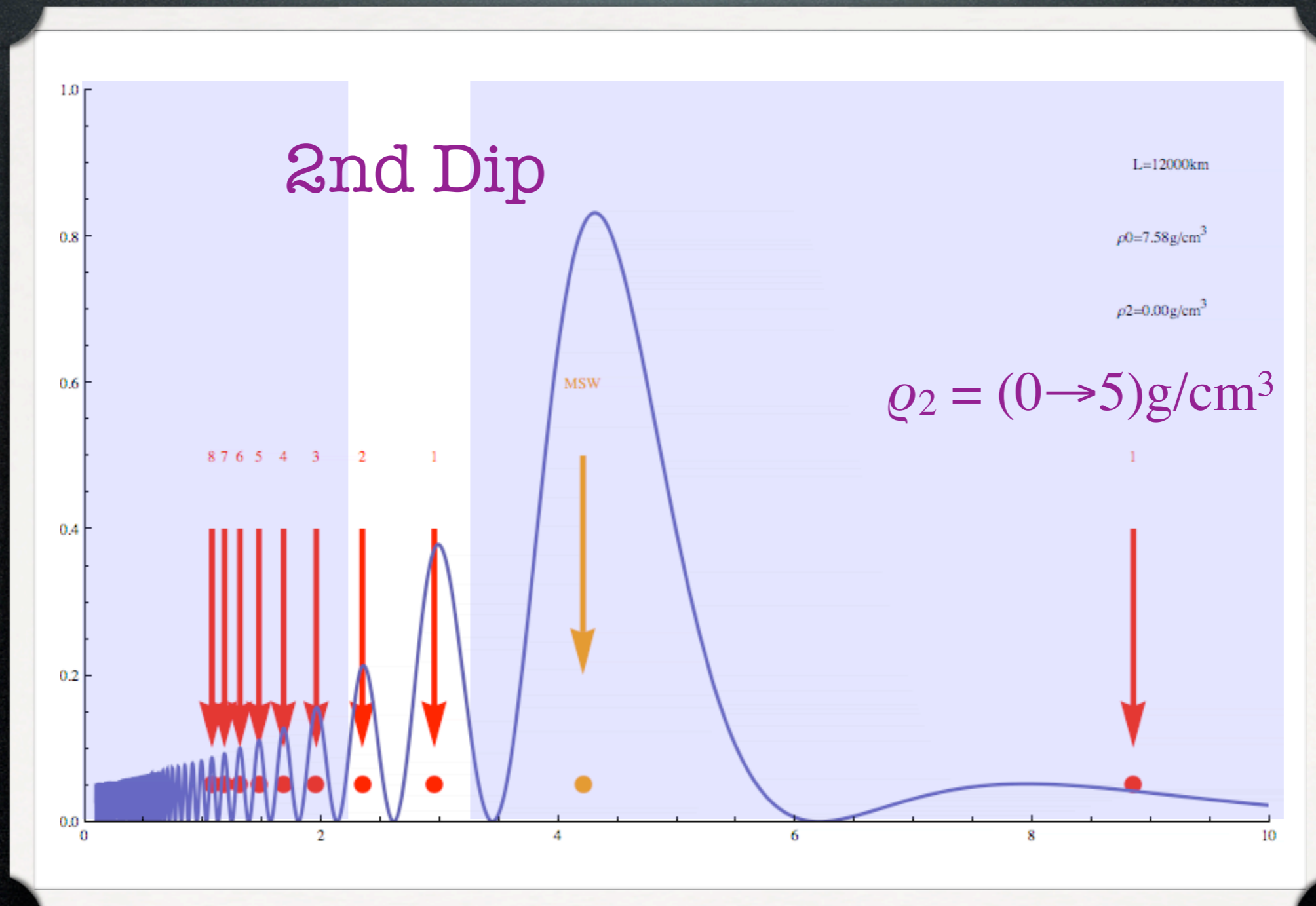


# Effect of the Mode 1





# Effect of the Mode 2



# Beyond a single swing, $\xi > 1$

$$z''(\xi) + (\omega_0^2 + \alpha_n \cos 2n\pi\xi - i\beta_n \sin 2n\pi\xi + \gamma_n \cos 4n\pi\xi) z(\xi) = 0$$

- What if the matter profile is repetitively applied?

Approximate solution at  $\omega_0 = n\pi$ :

$$z(\xi) = \frac{i\Delta \sin 2\theta}{2\Omega(4n^2\pi^2 - \alpha_n - 2\Omega^2)} \left[ -i\beta_n \sin \Omega\xi \sin n\pi\xi + (\alpha_n + 2\Omega^2) \sin \Omega\xi \cos n\pi\xi - 4n\pi\Omega \cos \Omega\xi \sin n\pi\xi \right],$$

$$\text{where } \Omega = \pm\sqrt{2}n\pi \left[ 1 - \sqrt{1 - \frac{\beta_n^2 - \alpha_n^2}{(2n\pi)^4}} \right]^{1/2}$$

$$|z(\xi)|^2 = \frac{\Delta^2 \sin^2 2\theta}{4\Omega^2(4n^2\pi^2 - b - 2\Omega^2)^2} \times$$

$$\left[ (b + 2\Omega^2)^2 \sin^2 \Omega\xi \cos^2 n\pi\xi + (16n^2\pi^2\Omega^2 \cos^2 \Omega\xi + c^2 \sin^2 \Omega\xi) \sin^2 n\pi\xi \right. \\ \left. - 8n\pi\Omega(b + 2\Omega^2) \cos \Omega\xi \sin \Omega\xi \cos n\pi\xi \sin n\pi\xi \right]$$

# Beyond a single swing, $\xi > 1$

$$z''(\xi) + (\omega_0^2 + \alpha_n \cos 2n\pi\xi - i\beta_n \sin 2n\pi\xi + \gamma_n \cos 4n\pi\xi) z(\xi) = 0$$

$$|z(\xi)|^2 = \frac{\Delta^2 \sin^2 2\theta}{4\Omega^2(4n^2\pi^2 - b - 2\Omega^2)}$$

$$\left[ (b + 2\Omega^2)^2 \sin^2 \Omega\xi \cos^2 n\pi - 8n\pi \Omega \cos \Omega\xi \sin 2n\pi \right]$$

$$\Omega = \pm \sqrt{2}n\pi \left[ 1 - \frac{b + 2\Omega^2 \cos \Omega\xi \sin 2n\pi}{4n^2\pi^2 - b - 2\Omega^2} \right]$$

Mode 1

$$\rho_1 = -0.3 \text{ g/cm}^3$$

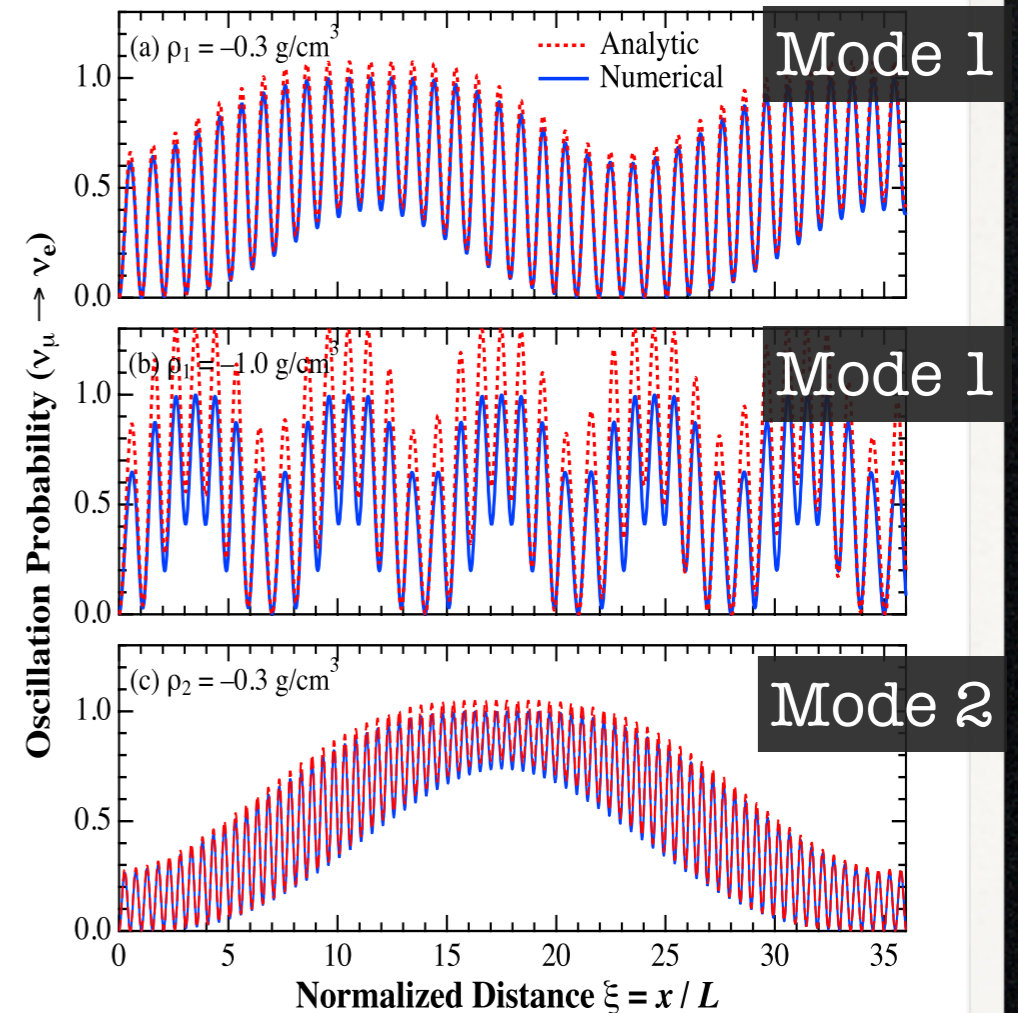
Mode 1

$$\rho_1 = -1.0 \text{ g/cm}^3$$

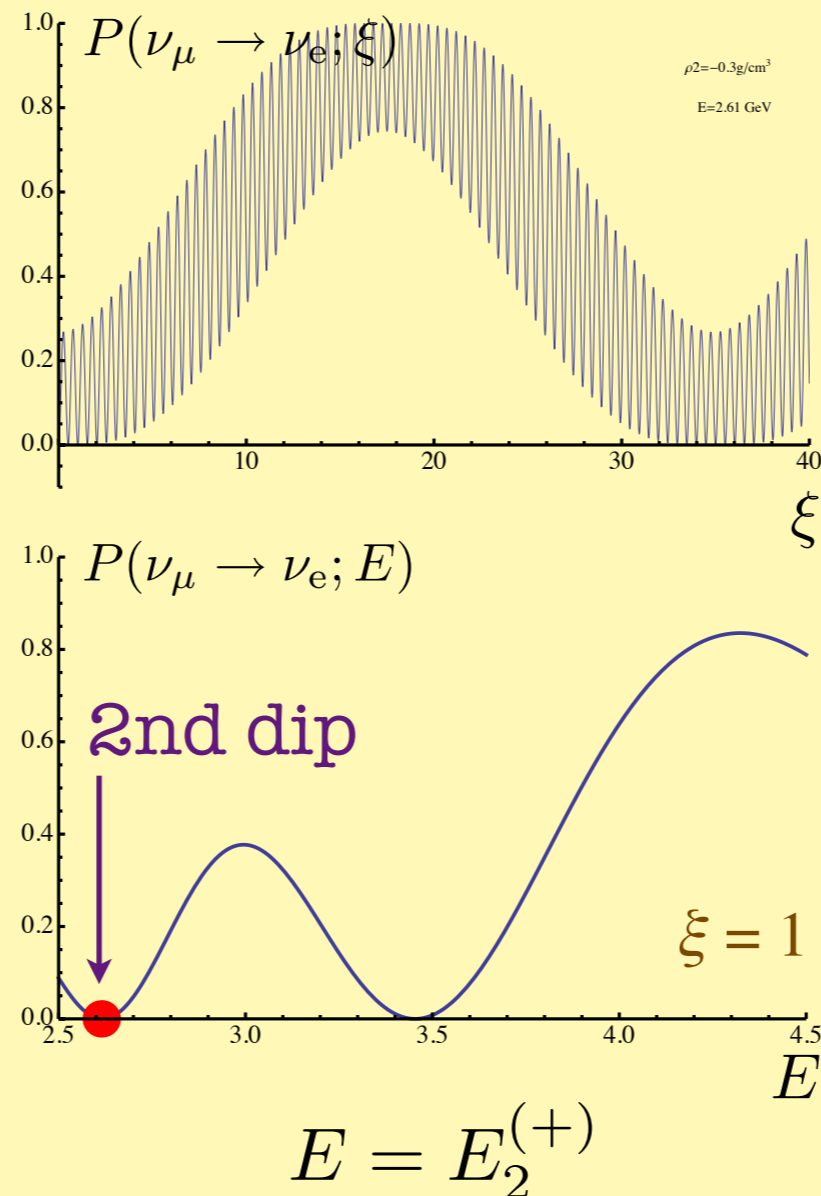
Mode 2

$$\rho_2 = -0.3 \text{ g/cm}^3$$

- Resonant growth of the appearance
- Resonance
- Unitarity ensured



# Off the Resonance Energy



# Off the Resonance

- Off-the-resonance ( $\omega_0 \neq n\pi$ , but  $\omega_0 \simeq n\pi$ )

$$z(\xi) = \frac{\Delta \sin 2\theta}{2\Omega[4n^2\pi^2 - \alpha_n - 2\Omega^2 + 2(\omega_0^2 - n^2\pi^2)]} \times$$

$$\left\{ i[\alpha_n + 2\Omega^2 - 2(\omega_0^2 - n^2\pi^2)] \sin \Omega\xi \cos n\pi\xi + \beta_n \sin \Omega\xi \sin n\pi\xi - i4n\pi\Omega \cos \Omega\xi \sin n\pi\xi \right\}$$

where

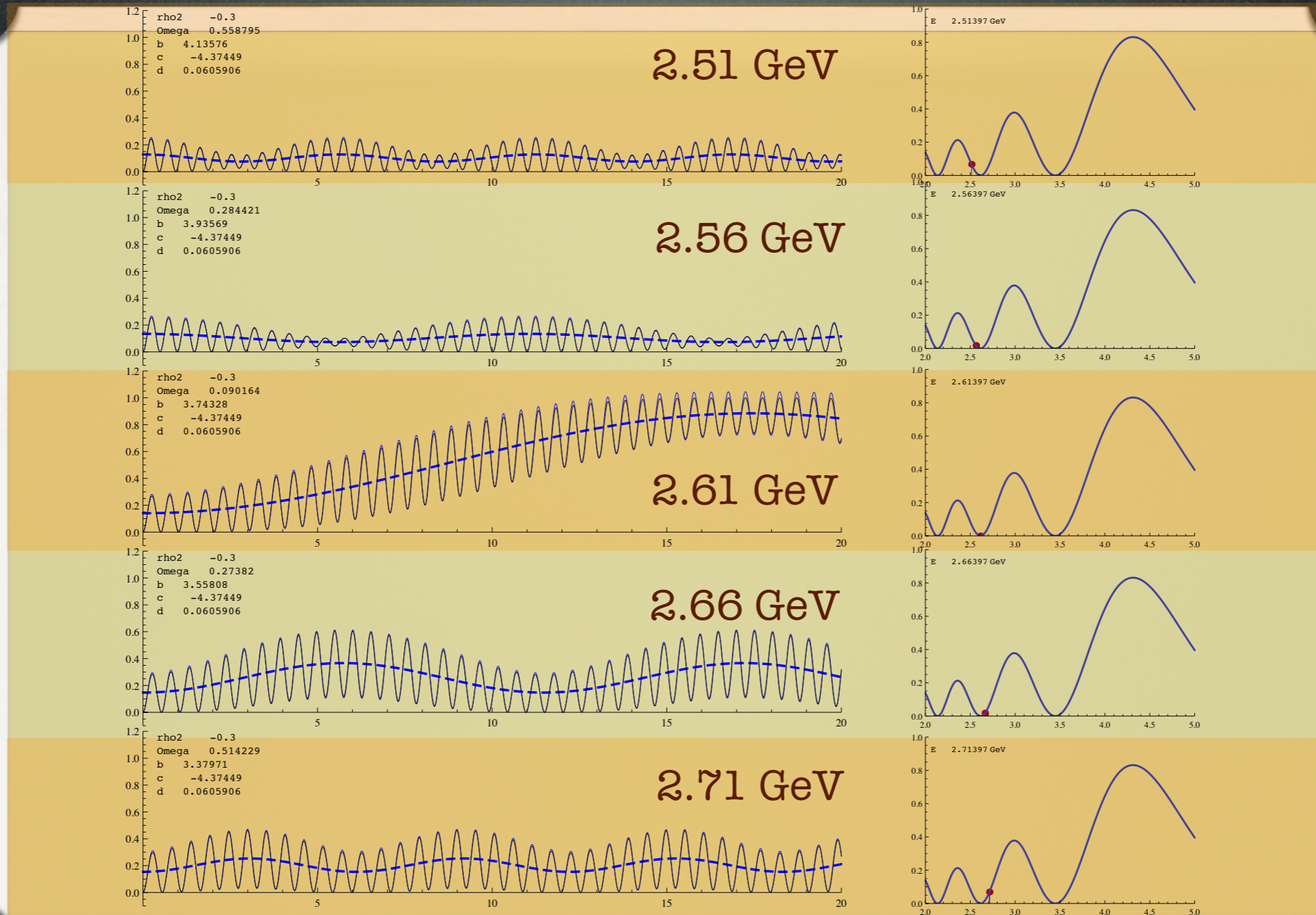
$$\Omega \equiv \sqrt{2n\pi} \left[ 1 + \frac{\omega_0^2 - n^2\pi^2}{2n^2\pi^2} - \sqrt{1 - \frac{\beta_n^2 - \alpha_n^2}{16n^4\pi^4} + \frac{\omega_0^2 - n^2\pi^2}{n^2\pi^2}} \right]^{1/2} \simeq \frac{\sqrt{\beta_n^2 - \alpha_n^2}}{4n\pi} \left[ 1 + \frac{2(\omega_0^2 - n^2\pi^2)^2}{\beta_n^2 - \alpha_n^2} - \frac{\omega_0^2 - n^2\pi^2}{4n^2\pi^2} + \frac{\beta_n^2 - \alpha_n^2}{128n^4\pi^4} \right]^{1/2}$$

$$|z(\xi)|^2 = \frac{\Delta^2 \sin^2 2\theta}{4\Omega^2 [4n^2\pi^2 - \alpha_n - 2\Omega^2 + 2(\omega_0^2 - n^2\pi^2)]^2} \times$$

$$\left\{ \left[ \alpha_n + 2\Omega^2 - 2(\omega_0^2 - n^2\pi^2) \right]^2 \sin^2 \Omega\xi \cos^2 n\pi\xi + (16n^2\pi^2\Omega^2 \cos^2 \Omega\xi + \beta_n^2 \sin^2 \Omega\xi) \sin^2 n\pi\xi \right.$$

$$\left. - 8n\pi\Omega [\alpha_n + 2\Omega^2 - 2(\omega_0^2 - n^2\pi^2)] \cos \Omega\xi \sin \Omega\xi \cos n\pi\xi \sin n\pi\xi \right\}$$

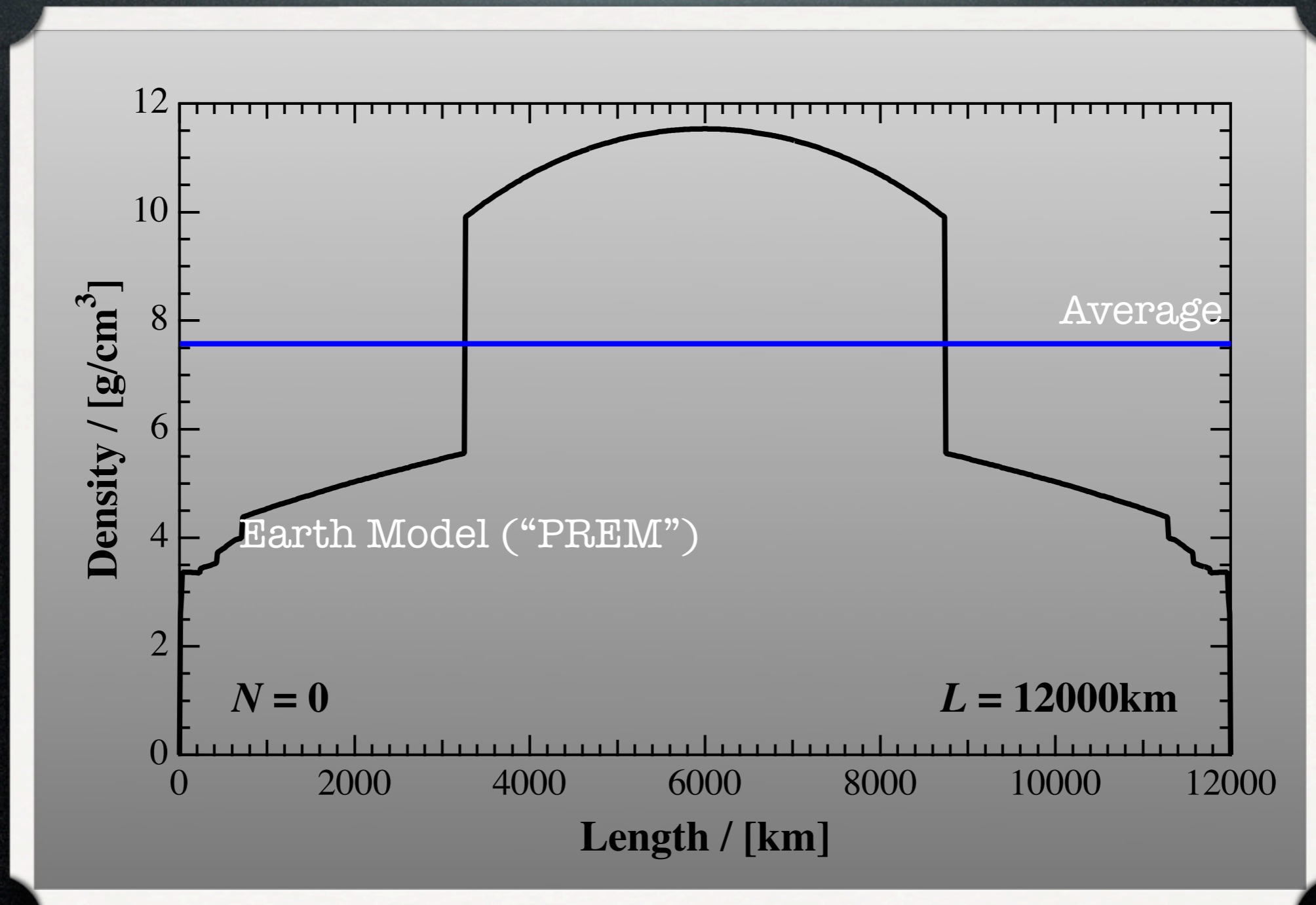
# Off the Resonance



# Summary & Outlook

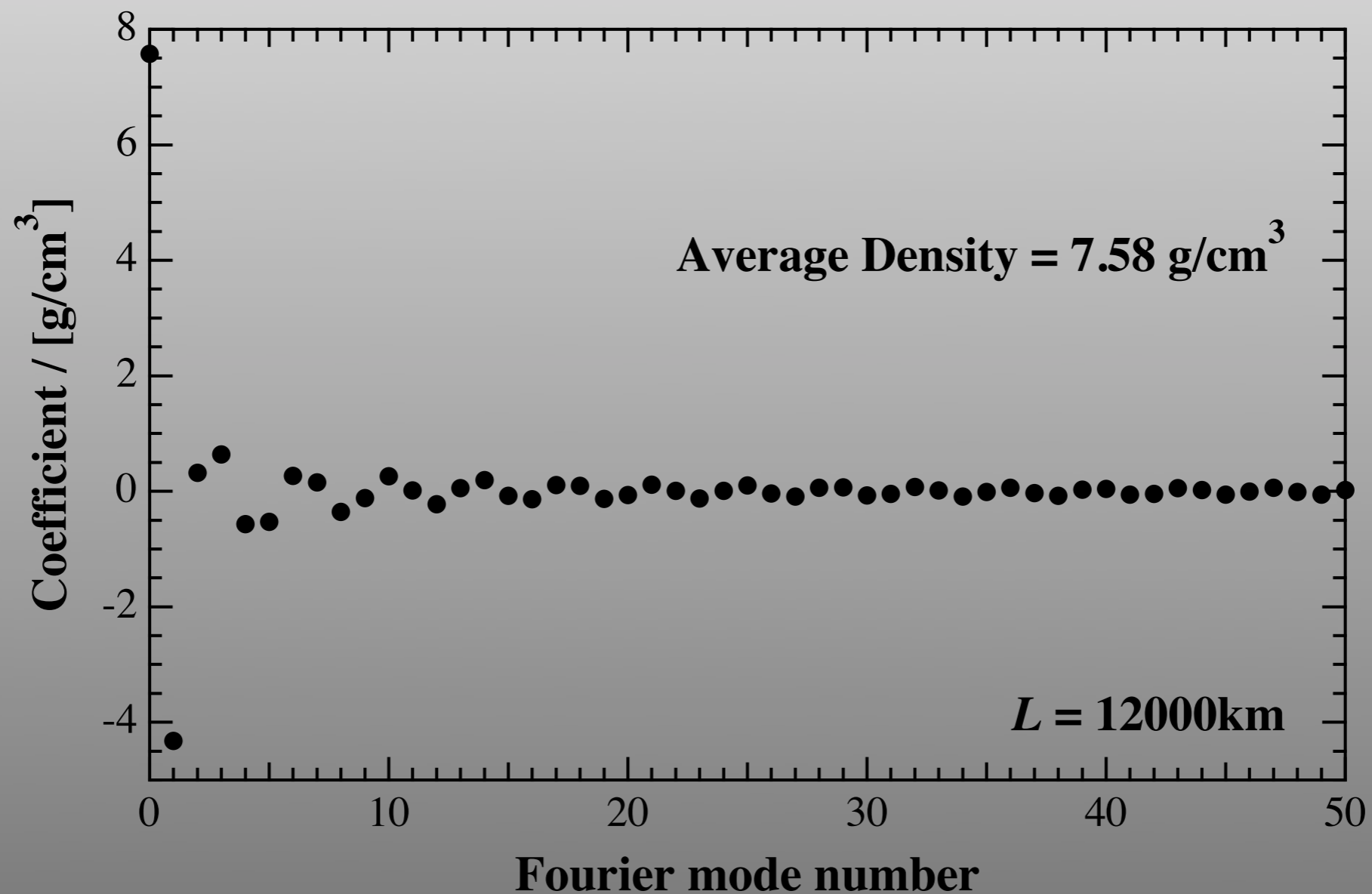
- **Fourier analysis** is shown to be powerful to account for the parametric resonance in neutrino oscillation.
  - Neutrino oscillation under inhomogeneous matter is analyzed, analytically and numerically.
  - **n-th Fourier mode of the matter profile affects at around the n-th dip of the appearance probability.**
  - Inhomogeneity effect leads to parametric resonance, resulting in a slow oscillation.
  - Resonance effect diminishes as the energy goes off the resonance point.
- Search for observable effects of interst (at  $\xi = 1$ )
  - Beyond a small potential: Large effect in a single swing?
- Further mathematical investigation

# Matter density on a baseline

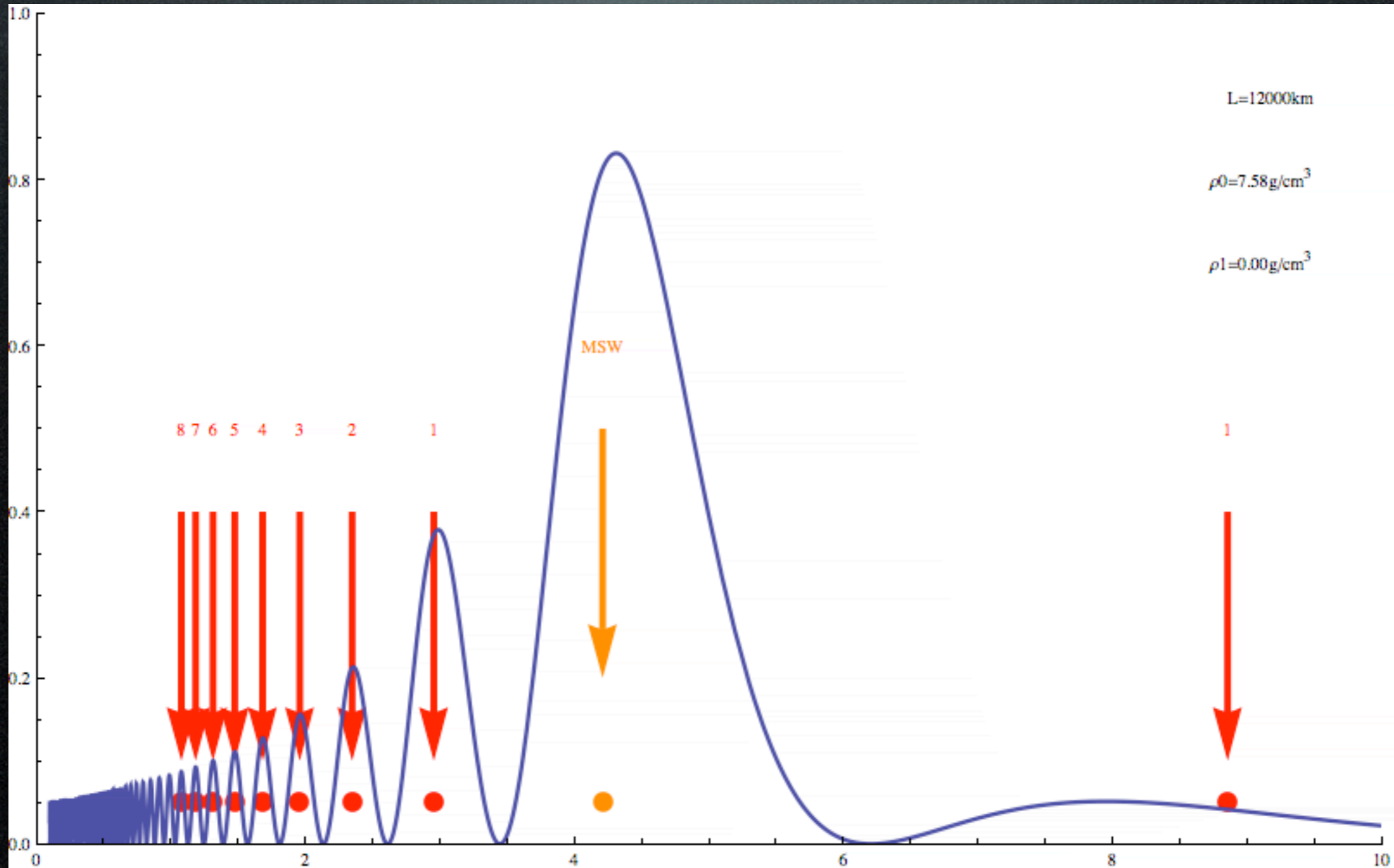




# Fourier coefficients



# First-mode effect



# Second-mode effect

